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The Broad Fe K α Line in MCG-6-30-15

J. Wilms¹ *, C. S. Reynolds², M. C. Begelman³, E. Kendziorra¹, R. Staubert¹, J. Reeves⁴ and S. Molendi⁵

- ¹ IAA Tübingen, Astronomie, Sand 1, 72076 Tübingen, Germany
- ² Dept. of Astronomy, University of Maryland, College Park, MD 20742, USA
- ³ JILA, University of Colorado, Boulder, CO 80309, USA
- ⁴ Laboratory for High Energy Astrophysics, NASA Goddard Space Flight Center, Greenbelt, MD 20771, USA
- ⁵ ASF CNR, Via Bassini 15, 20133 Milano, Italy

Abstract. Fluorescent iron $K\alpha$ lines in active galactic nuclei serve as probes for the behavior of material under the extreme conditions close to the supermassive black hole. We discuss the evidence for such relativistic effects in observational data from Seyfert galaxies taken with the EPIC cameras on board XMM-Newton.

Key words. black hole physics - galaxies: active - accretion, accretion disks

1. Introduction

One of the "holy grails" of modern astrophysics is the search for observational evidence of the existence of black holes (BHs). A first step towards finding this "holy grail" would be the proof that the behavior of material close to large mass concentrations is consistent with what is predicted by general relativity. X-ray observations of Fe K α lines from active galactic nuclei (AGN) are currently one of the most likely means to achieve this goal. In this contribution we first describe how broad Fe K α lines are produced (§2) and then give an overview of our results from *XMM-Newton* guaranteed time observations of the broad Fe K α line source MCG-6-30-15 (§3).

2. Relativistic Fe K α Lines

In the standard model of AGN, $1-2 M_{\odot} \text{ yr}^{-1}$ accrete via an accretion disk onto a central supermassive BH (Antonucci, 1993, and therein). The detailed structure of the accretion flow is unknown. While the disk has likely a thermal spectrum - which is observed as the "big blue bump" in the UV and the soft X-rays – X-ray data show a further hard power law component that turns over into an exponential cutoff at ~100 keV. Current thought is that this hard component is due to Comptonization of soft disk photons in a hot electron gas - the "accretion disk corona" (ADC) - which surrounds the disk (e.g., Haardt et al., 1997, see also Fig. 1). A possible origin of the ADC are the MHD instabilities thought to be responsible for the viscosity of the accreting material (Balbus & Hawley, 1998). Efficient Compton cooling of the electrons prevents the formation of an ADC which fully covers the accre-

Send offprint requests to: J. Wilms

^{*} present address: Department of Physics, University of Warwick, Coventry, CV4 7AL, U.K.

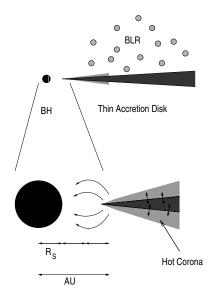


Fig. 1. Sketch of the inner region of an AGN accretion disk.

tion disk, so that patchy coronae or geometries where the disk and the ADC are separated are more likely (e.g., Dove et al., 1997).

Regardless of the geometry, a fraction of the hard ADC photons are expected to be intercepted by the disk. These photons are either downscattered by electrons in the accretion disk, or absorbed by the accretion disk material. The detailed physics of this process strongly depends on the ionization state of the accretion disk (see, e.g., Ballantyne et al. 2002 and Nayakshin & Kallman 2001), however, the X-ray spectrum from a disk irradiated by a power-law is roughly characterized by a Compton reflection hump peaking at ~30 keV (Lightman & White, 1988). Below ~10 keV, this "Compton reflection hump" is dominated by absorption and fluorescent line emission. The most prominent of these lines is the Fe K α line at \sim 6.4 keV, which provides a very distinct spectral feature.

Since the reprocessing features (Fe $K\alpha$ line and reflection hump) come from close to the BH, distant observers are able to see the influence of the large speeds (the Keplerian speed at the innermost stable circular orbit is a significant fraction of the speed of light) and the effects of strong gravity in the observed X-ray

spectrum. As a result of these effects, the observed Fe $K\alpha$ line shape is strongly broadened and skew-symmetric. As is shown elsewhere (e.g., Reynolds & Nowak, 2003; Wilms et al., 1998, and therein), the study of the Fe $K\alpha$ line profile allows in principle the determination of the BH's angular momentum and of the Fe $K\alpha$ emissivity.

The first secure detection of a relativistic Fe K α line came from a long ASCA observation of the Seyfert galaxy MCG-6-30-15 (Tanaka et al., 1995). Later observations showed that broadened Fe K α lines are common in the X-ray spectra of Seyfert galaxies (e.g., Nandra et al., 1997; Molendi et al., 1997; Lee et al., 1998). With the Chandra HETGS and the XMM-Newton EPIC, these earlier results have been confirmed and extended, e.g., with new detections of broad lines in sources such as the Seyfert 1.9 galaxy MCG-5-23-16 (Dewangan et al., 2003) or the quasar O0056-363 (Porquet & Reeves, 2003). At the same time, some AGN thought to exhibit broad lines were shown to have more complex spectra (e.g., Reeves, 2003).

XMM-Newton Observations of MCG-6-30-15

With the launch of XMM-Newton in 1999, an instrument allowing one to measure the properties of Fe K α lines with ~150 eV resolution and a high signal-to-noise ratio has become available. It was thus natural to use the XMM-Newton guaranteed time (GT) to further our understanding of MCG-6-30-15 with a 100 ksec long observation of the source in 2000 June. Luckily enough, the XMM-Newton GT observation caught the AGN during a time of low flux, where the Fe K α line was expected to be broad (Iwasawa et al., 1996). Here, we summarize our Fe K α line results from this observation, see Wilms et al. (2001) and Reynolds et al. (2004) for more detailed information. We note that a second and longer XMM-Newton observation of MCG-6-03-15 was performed as part of XMM-Newton's open time program in 2001 July and August, when the source was 70% brighter than in the GT observation. The main results of the

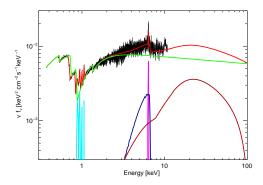


Fig. 2. Inferred vf_v spectrum of MCG-6-30-15 from the EPIC pn data, showing the model components of the best fit model. Note that the data have been only slightly rebinned.

GT observation are confirmed by these data which have been analyzed in a series of papers by Fabian and collaborators (e.g., Fabian & Vaughan, 2003; Fabian et al., 2002, and therein).

In order to determine the Fe K α parameters, an understanding of the underlying continuum spectrum is required. This determination of the continuum is complicated by features in the soft X-rays, which are thought to be either absorption edges caused by ionized material along the line of sight – the "dusty warm absorber" (Lee et al., 2001, and references therein) - or relativistically broadened O and Ni recombination lines (Branduardi-Raymont et al., 2001; Sako et al., 2003). The effect of these features is measurable at energies up to 2 keV. To gauge the effect the soft X-rays have on the modeling of the Fe K α line, we model both the 2-10 keV and the 0.5-10 keV EPICpn data. Furthermore, in modeling the 0.5-10 keV data we gauge the stability of our interpretation of the > 2 keV spectrum by using both, an empirical warm absorber model and emission lines for describing the soft X-rays. The qualitative results of our data modeling are the same in both cases (Reynolds et al., 2004).

The continuum >2 keV is modeled by a power law, ionized reflection, and the Fe $K\alpha$ complex. The data clearly show that the Fe $K\alpha$ line is the superposition of a narrow feature at 6.4 keV, which has constant flux dur-

ing the observation and likely originates from material far away from the central source, as well as the relativistic broadened feature. To describe the relativistic smearing of the Fe K α line and the continuum we first assume a power law emissivity per disk unit area of the form $\epsilon \propto r^{-\beta}$ for the line. Fitting the data shows that $\beta > 3$, i.e., the emissivity of the Fe K α line is strongly concentrated towards the inner rim of the accretion disk (Wilms et al., 2001) such that the observed Fe K α line is strongly broadened (Fig. 2). Since simple ADC models predict that the Fe K α line emissivity is proportional to the energy dissipated in the accretion disk, one would expect β < 3. This led us to the conclusion that a further energy dissipating process must be at work close the the central compact object, e.g., magnetic coupling between a Kerr BH and the accretion disk, possibly by torquing the inner edge of the disk (Agol & Krolik, 2000; Li, 2002). Based on such models, we have reanalyzed our GT data with more physical models, confirming our earlier conclusions and showing that the Fe K α line can be well described by emission from a strongly torqued accretion disk which emits most X-rays from within a few gravitational radii (Reynolds et al., 2004).

4. Conclusion

We have described the results of the analysis of XMM-Newton guaranteed time data of MCG-6-30-15, showing that the instruments on XMM-Newton are well suited for the study of the behavior of material in the strong gravitational field close to a BH. Clearly, future theoretical and observational work is needed since XMM-Newton and other satellites have proven that MCG-6-30-15 is strongly variable. The study of this variability will improve our ability to understand the source, by studying correlations between the different spectral components. Furthermore, alternative models for the line broadening have to be investigated. At the moment, the only other model explaining the line shape is the "lamppost model" (e.g., Martocchia et al., 2002), where the X-ray source is assumed to be close to the spin axis of the black hole and above the accretion disk.

Changes in the height of such an X-ray source can explain the correlation between the continuum intensity and the strength of the X-ray reflection feature (Fabian & Vaughan, 2003). Further high signal to noise data of MCG-6-30-15, however, both from XMM-Newton and Chandra, as well as from future missions such as XEUS or Constellation-X, are required to disentangle the spectral components in this complicated source and shed further light on the reality of relativistic effects in the universe.

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References

Agol, E. & Krolik, J. H. 2000, ApJ, 528, 161 Antonucci, R. 1993, ARA&A, 31, 473

Balbus, S. A. & Hawley, J. F. 1998, Rev. Mod. Phys., 70, 1

Ballantyne, D. R., Ross, R. R., & Fabian, A. C. 2002, MNRAS, 336, 867

Branduardi-Raymont, G., et al. 2001, A&A, 365, L140

Dewangan, G. C., Griffiths, R. E., & Schurch, N. J. 2003, ApJ, 592, 52

Dove, J. B., et al. 1997, ApJ, 487, 759

Fabian, A. C. & Vaughan, S. 2003, MNRAS, 340, L28

Fabian, A. C., et al. 2002, MNRAS, 335, L1Haardt, F., Maraschi, L., & Ghisellini, G. 1997, ApJ, 476, 620

Iwasawa, K., et al. 1996, MNRAS, 282, 1038 Lee, J. C., et al. 1998, MNRAS, 300, 583 Lee, J. C., et al. 2001, ApJ, 554, L13 Li, L.-X. 2002, Phys. Rev. D, 65, 84047

Lightman, A. P. & White, T. R. 1988, ApJ, 335, 57

Martocchia, A., Matt, G., & Karas, V. 2002, A&A, 383, L23

Molendi, S., et al. 1997, A&A

Nandra, K., et al. 1997, ApJ, 477, 602

Nayakshin, S. & Kallman, T. R. 2001, ApJ, 546, 406

Porquet, D. & Reeves, J. N. 2003, A&A, 408, 119

Reeves, J. 2003, in Active Galactic Nuclei: from Central Engine to Host Galaxy,, ed. S. Collin, et al. in press (astro-ph/0211381)

Reynolds, C. S. & Nowak, M. A. 2003, Phys. Rep., 377, 389

Reynolds, C. S., et al. 2004, MNRAS, submitted

Sako, M., et al. 2003, ApJ, in press (astroph/0112436)

Tanaka, Y., et al. 1995, Nature, 375, 659

Wilms, J., et al. 2001, MNRAS, 328, L27

Wilms, J., Speith, R., & Reynolds, C. S. 1998, in Black Holes: Theory and Observation, ed. F. W. Hehl, et al., LNP 514 (Berlin, Heidelberg: Springer), 69